## Incommensurate magnetic structure in Y<sub>(1-x)</sub>Tb<sub>x</sub>Mn<sub>6</sub>Sn<sub>6</sub>.

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We have investigated the magnetic structure in the intermetallic compounds  $Tb_xY_{1-x}Mn_6Sn_6$  ( $0 \le x \le 0.25$ ) by means of small angle neutron scattering (SANS). In such intermetallic materials the competition between different exchange interactions results in a variety of the magnetic structures and properties. Thus, giant magnetoresistance was recently found in this type of materials [1].

The YMn<sub>6</sub>Sn<sub>6</sub> is a compound crystallized in hexagonal HfFe<sub>6</sub>Ge<sub>6</sub>-type structure (space group P6/mmm) [2]. The lattice has an intrinsically layered structure, where Mn atoms are organized in a so-called "kagome" lattice within the ab planes, which are stacked along the c axis with Y and Sn<sub>3</sub> atomic planes between them. The Mn-Mn interplane distance through Sn<sub>3</sub> atomic plane is slightly larger than that through Y atomic plane. Below the Neel temperature  $T_N$ =333K the YMn<sub>6</sub>Sn<sub>6</sub> compound has an incommensurate periodic structure [3]. It is believed that spins are ordered in the helix along the c axis similar to a simple planar helimagnet [4].

To make clear a role of different exchange interactions we studied samples where Y is partially substituted by Tb. The substitution of magnetic Tb for Y in the  $Tb_xY_{1-x}Mn_6Sn_6$  compounds changes the magnetic ordering type from incommensurate antiferromagnetic to ferrimagnetic at the concentration  $x \approx 0.2$ . The ferrimagnets at x > 0.2 show a change of the magnetocrystalline anisotropy with increase of temperature from an easy-axis type to an easy-plane type through a conical phase [5].

The use of SANS allows one to observe large-scale magnetic ordering modes. Two scattering contributions are well distinguished. The first one is a diffraction peak, which originates from the scattering on the magnetic long periodic structure. The second one is a small angle scattering attributed to the magnetic inhomogeneities such as domains or critical spin fluctuations. The period of the magnetic structure calculated from diffraction peak position for T= 175K is  $d_{TB0}$ =3.76 nm for YMn<sub>6</sub>Sn<sub>6</sub> and is  $d_{TB.2}$ =7.06 nm for Tb<sub>0.2</sub>Y<sub>0.8</sub>Mn<sub>6</sub>Sn<sub>6</sub>. Coherent lengths of magnetic structure are w<sub>TB0</sub>=17 nm and w<sub>TB.2</sub>=28 nm. Experimental data suggest that the temperature behavior of the magnetic structure for YMn<sub>6</sub>Sn<sub>6</sub> is very different from that for the simple planar helimagnet of Ho or Dy- type [4]. The temperature evolution of the magnetic structure for samples with Tb- substitution is even more complex. The c-T phase diagram is obtained on the basis of SANS measurements.

## References

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